

Modelowanie zjawiska kondensacji kapilarnej w materiałach porowatych



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Capillary condensation in cylindrical pores: Monte Carlo study of the interplay of surface and finite size effects

A. Winkler, D. Wilms, P. Virnau, and K. Binder

Citation: *J. Chem. Phys.* **133**, 164702 (2010); doi: 10.1063/1.3502684

Capillary condensation of adsorbates in porous materials

Toshihide Horikawa, D.D. Do*, D. Nicholson

Advances in Colloid and Interface Science 169 (2011) 40–58

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ARTICLE

pubs.acs.org/JPCB

Effect of Pore Size on the Condensation/Evaporation Transition of Confined Water in Equilibrium with Saturated Bulk Water

Ivan Brovchenko* and Alla Oleinikova*

Physical Chemistry, Dortmund University of Technology, Otto-Hahn-Strasse 6, Dortmund, D-44227, Germany

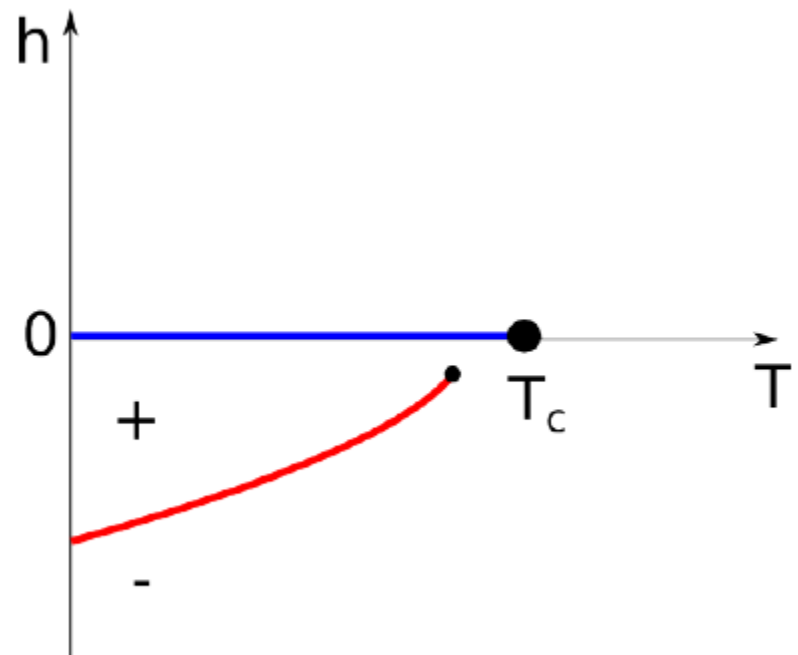
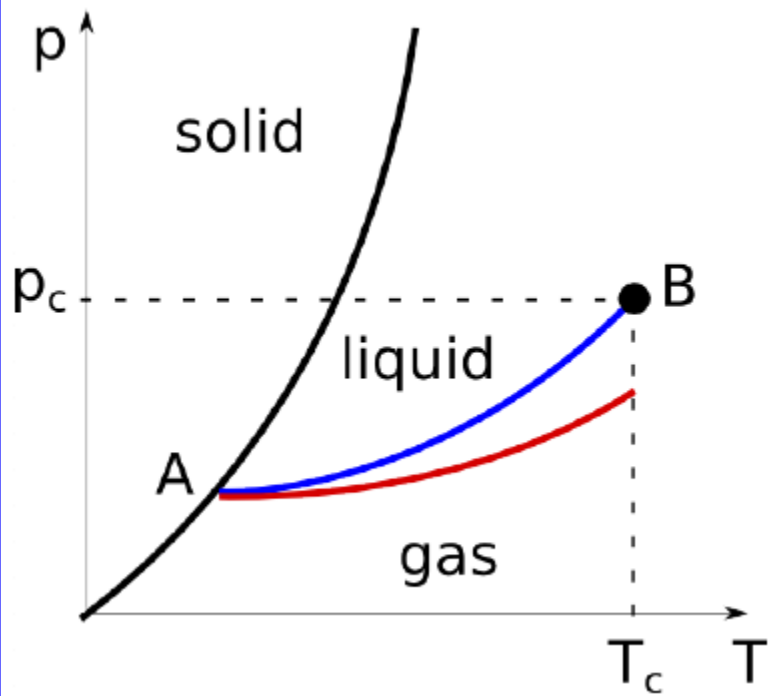
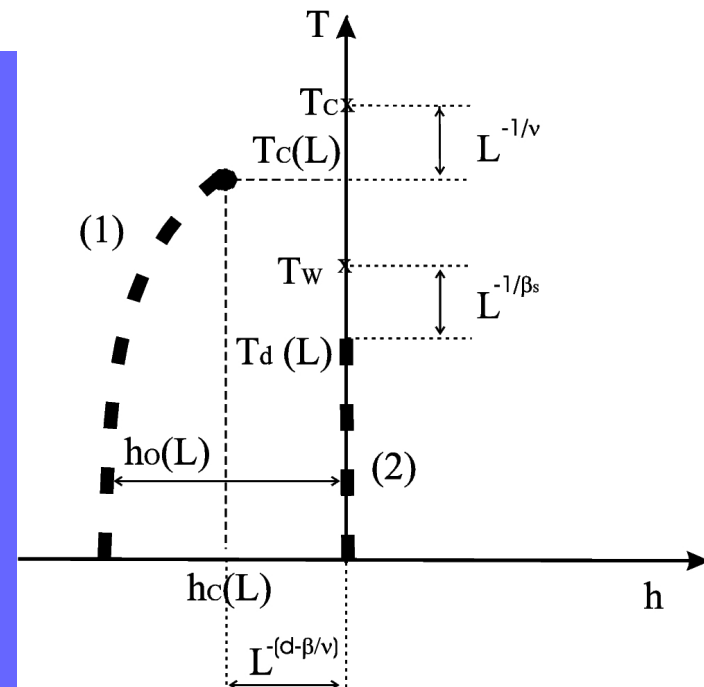
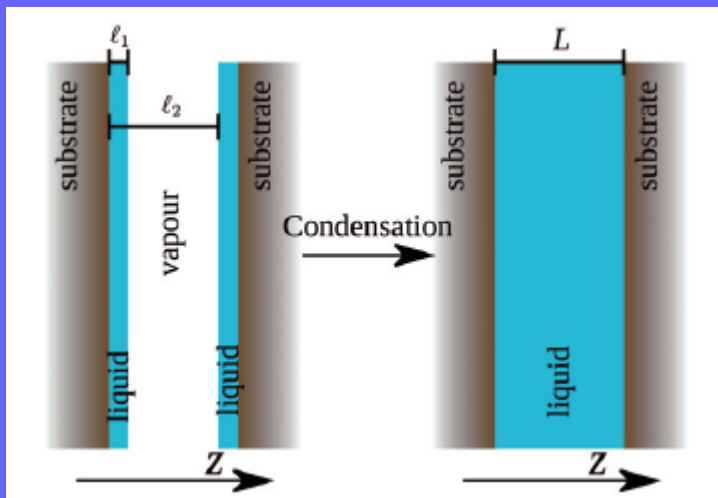


Diagram fazowy



Statistical mechanics of 2D classical spin systems on square lattice $N \times M$

Partition function Z is evaluated by means of the transfer matrix $T = \exp(-H / k_B T)$:

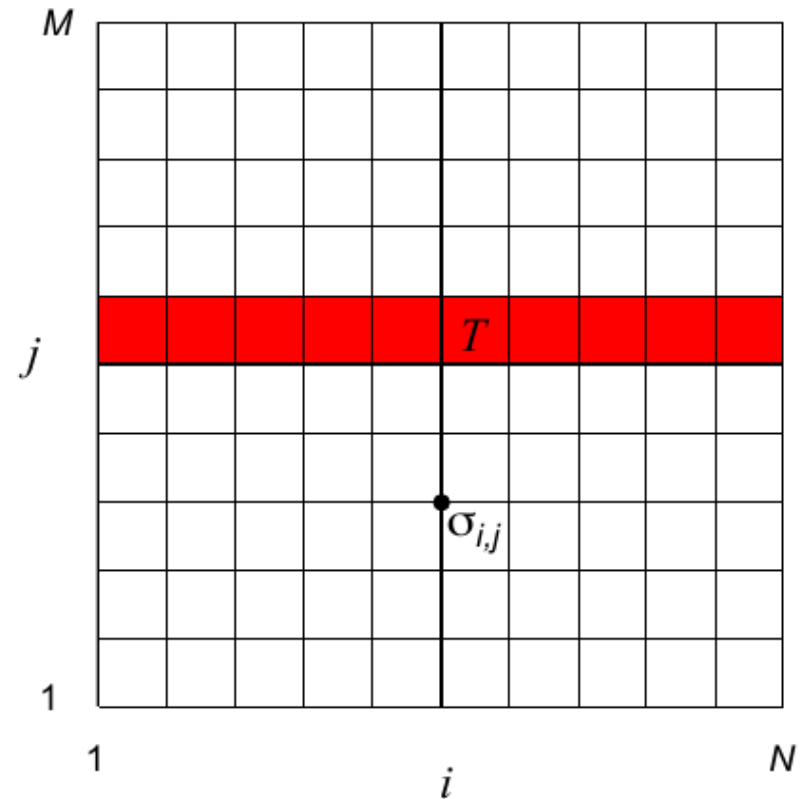
$$Z = \sum_{\{\sigma\}} \prod_{j=1}^M T(\sigma_{1,j} \cdots \sigma_{N,j} | \sigma_{1,j+1} \cdots \sigma_{N,j+1})$$

$$= \sum_{\{\sigma\}} \prod_{j=1}^M U^T \Lambda U$$

$$= \text{Tr} (U^T \Lambda U)^M = \text{Tr} \Lambda^M$$

$$= \sum_{k=1}^{2^N} \lambda_k^M$$

$$= \lambda_1^M \left[1 + \left(\frac{\lambda_2}{\lambda_1} \right)^M + \left(\frac{\lambda_3}{\lambda_1} \right)^M + \cdots + \left(\frac{\lambda_{2^N}}{\lambda_1} \right)^M \right]$$



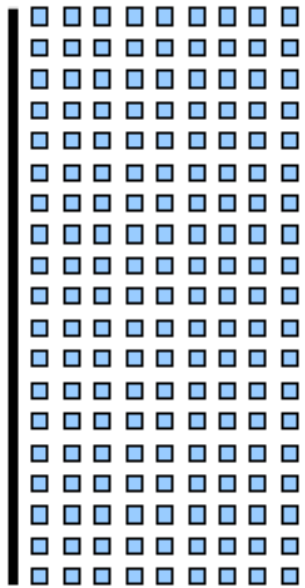
$$\lambda_1 > \lambda_2 \geq \lambda_3 \geq \cdots \geq \lambda_{2^N}$$

$$\begin{aligned}
f &= \lim_{M \rightarrow \infty} \frac{F}{M} = -k_B T \lim_{M \rightarrow \infty} M^{-1} \ln Z \\
&= -k_B T \ln \lambda_1 - \lim_{M \rightarrow \infty} \frac{k_B T}{M} \ln \left[1 + \underbrace{\left(\frac{\lambda_2}{\lambda_1} \right)^M + \left(\frac{\lambda_3}{\lambda_1} \right)^M + \dots + \left(\frac{\lambda_{2^N}}{\lambda_1} \right)^M}_{=0 \text{ if } \lambda_1 > \lambda_2 > \lambda_3 > \dots > \lambda_{2^N}} \right] \\
&= -k_B T \ln \lambda_1
\end{aligned}$$

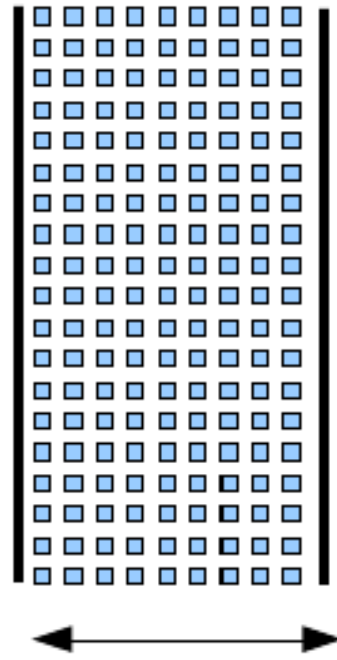
The free energy is proportional to the largest eigenvalue of the transfer matrix, and most of the thermodynamic functions can be evaluated

$$U \propto T^2 \frac{\partial f / T}{\partial T} \quad M \propto \frac{\partial f}{\partial H}$$

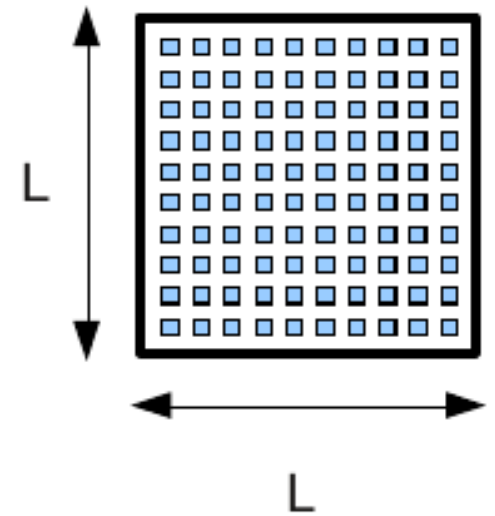
$$C \propto \frac{\partial U}{\partial T} \quad \chi \propto \frac{\partial f}{\partial H}$$



Semi-infinite
geometry



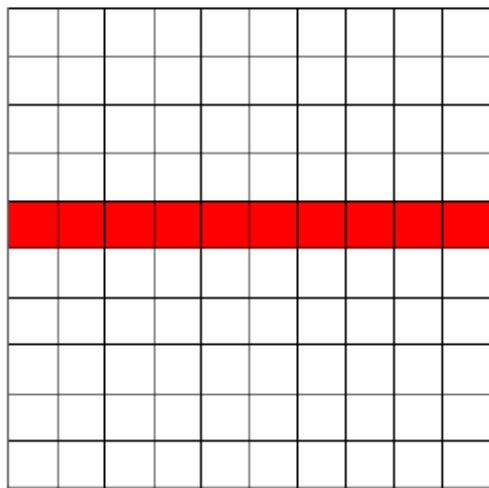
Strip-like
geometry



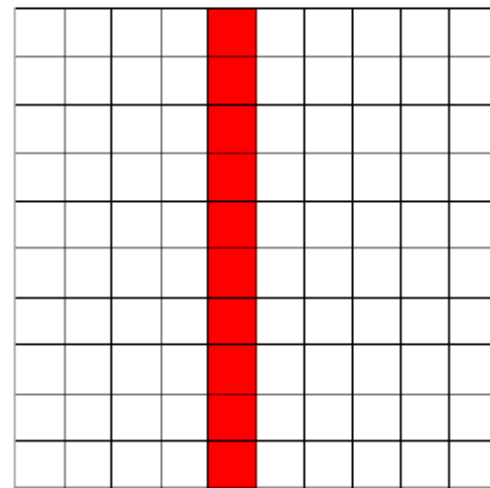
Square-like
geometry

There are many ways how to evaluate the partition function

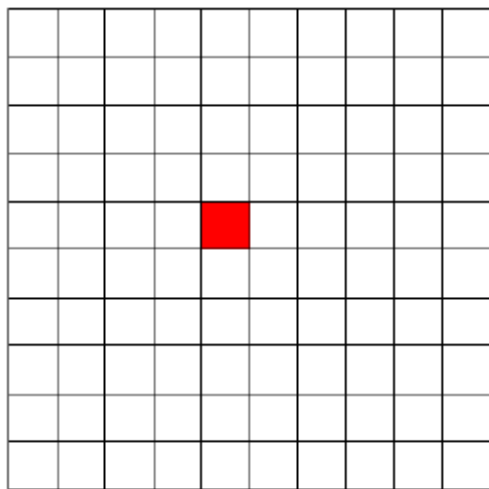
$$(1) \quad Z = \sum_{\{\sigma\}} \prod_{j=1}^M T(\sigma_{1,j} \cdots \sigma_{N,j} \mid \sigma_{1,j+1} \cdots \sigma_{N,j+1})$$



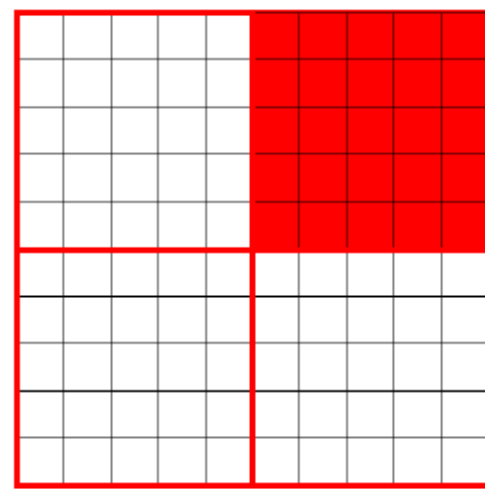
$$(2) \quad Z = \sum_{\{\sigma\}} \prod_{i=1}^N T(\sigma_{i,1} \cdots \sigma_{i,M} \mid \sigma_{i+1,1} \cdots \sigma_{i+1,M})$$



$$(3) \quad Z = \sum_{\{\sigma\}} \prod_{i=1}^N \prod_{j=1}^M W_B(\sigma_{i,j} \sigma_{i+1,j} \mid \sigma_{i,j+1} \sigma_{i+1,j+1})$$



$$(4) \quad Z = \sum_{\{\sigma\}} C\{\sigma\} \cdot C\{\sigma\} \cdot C\{\sigma\} \cdot C\{\sigma\}$$



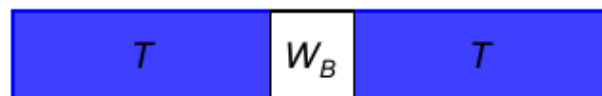
Simplified DMRG extension algorithm



$4 \times \infty$

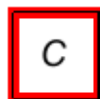


$6 \times \infty$

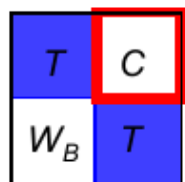


$8 \times \infty$

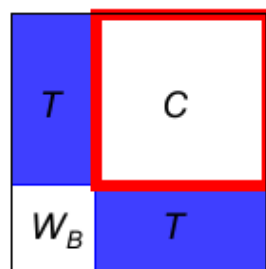
Simplified CTMRG extension algorithm



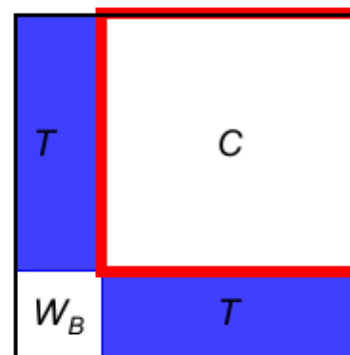
2×2



4×4

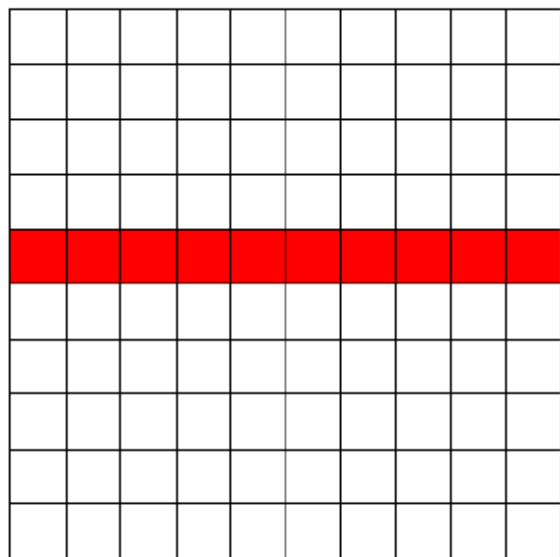


6×6

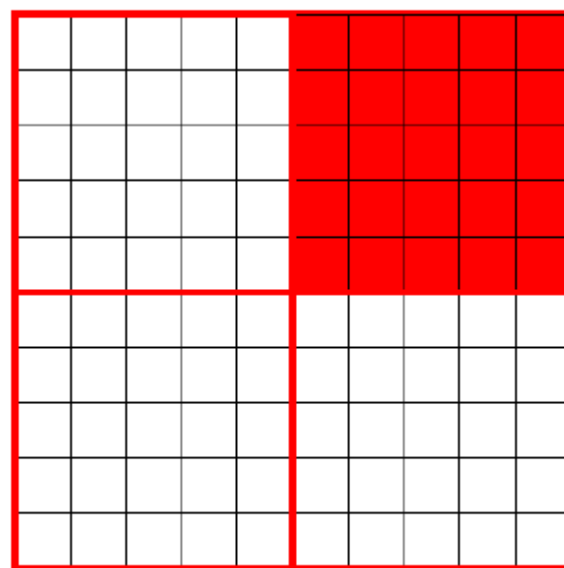


8×8

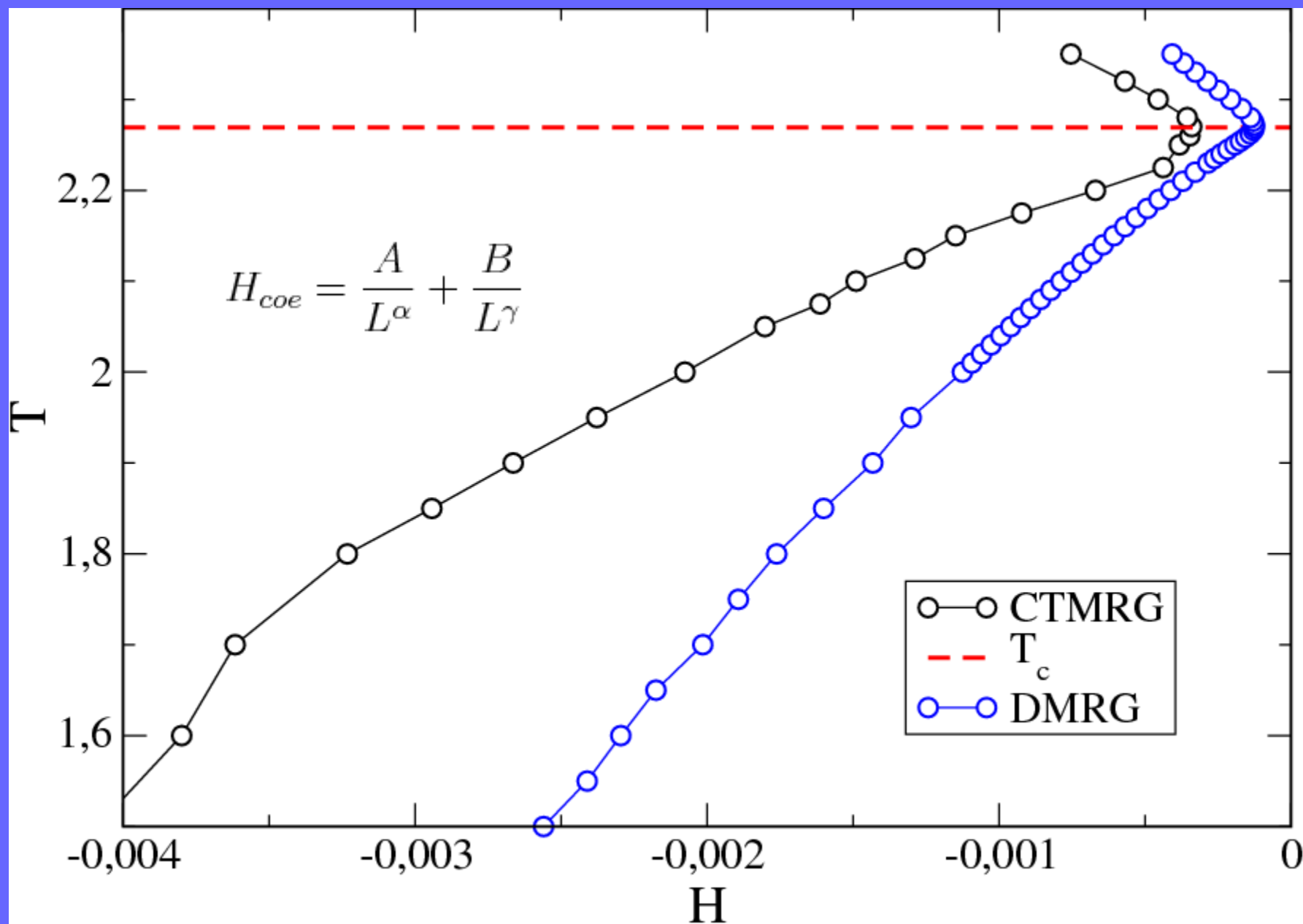
DMRG



CTMRG

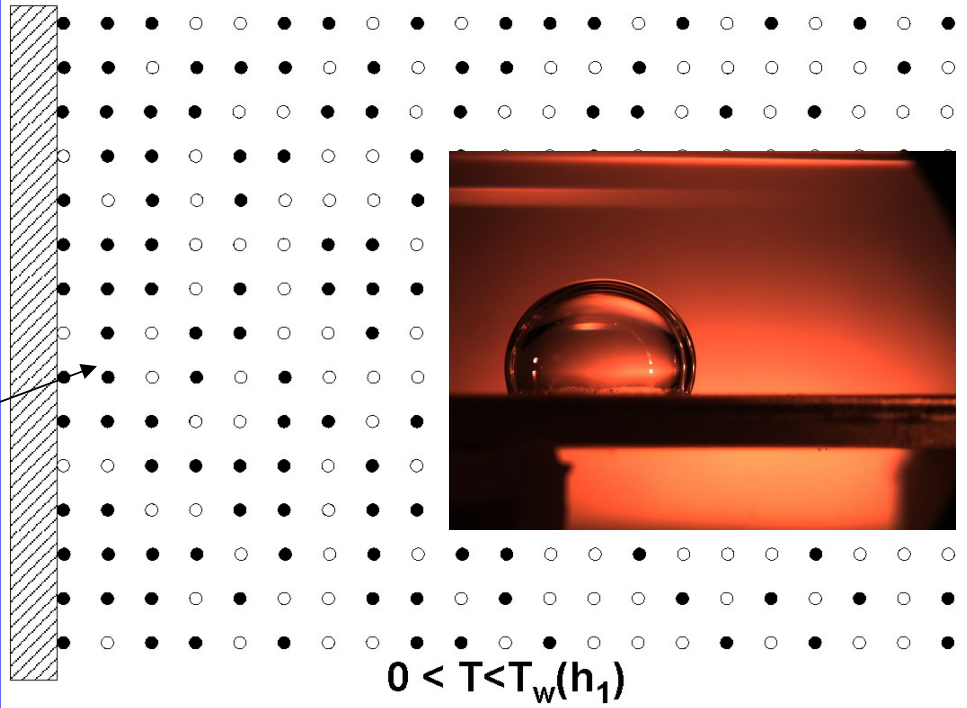


L = 500

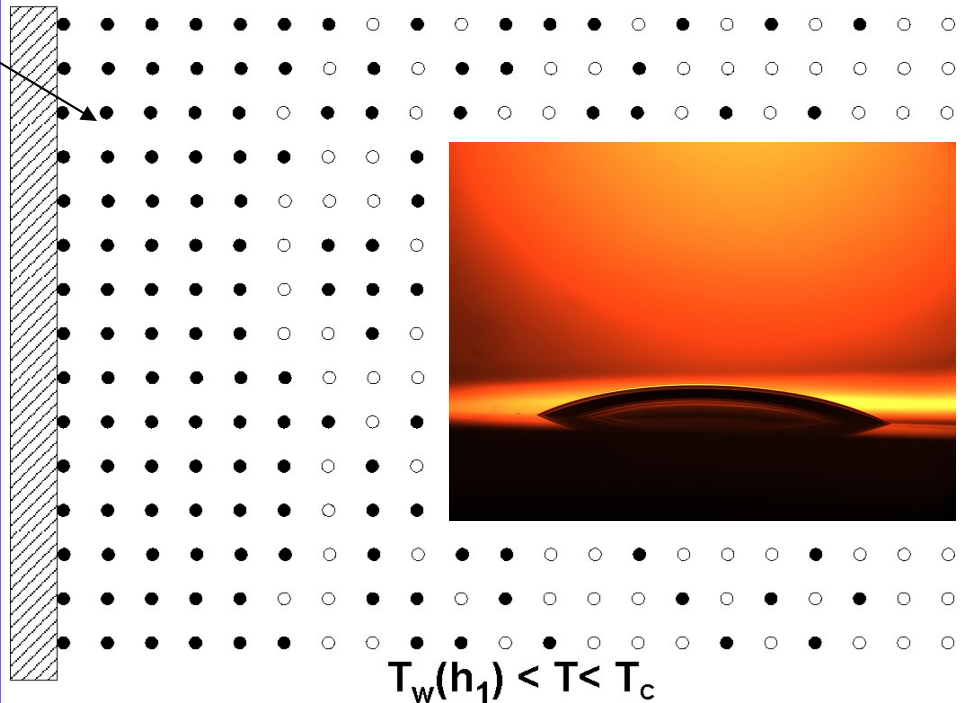


Zwilżanie częściowe

*W pobliżu linii
współistnienia !*



Zwilżanie pełne



Dziękuję za uwagę

